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1 **Methane production by CO<sub>2</sub>hydrogenation reaction with**  
2 **and without solid phase catalysis**

3  
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25 **Abstract**

26 Leveraging on the experience gained by designing and assembling a prototype experimental  
27 apparatus reusing CO<sub>2</sub> to produce methane in a reaction with H<sub>2</sub>, we are developing an  
28 alternative innovative cost effective and contaminant resistant synthetic strategy based on the  
29 replacement of the solid phase catalysis with a homogeneous gas phase process going through  
30 the forming CO<sub>2</sub><sup>2+</sup> molecular dications.

31 **Keywords:** methanation reactor; CO<sub>2</sub> waste; free methane; catalyst; plasma, dications.

32

33 **Introduction**

34 As reported during the 2<sup>nd</sup> SMARTCATs general meeting of the COST Action CM1404 we  
35 are carrying out a research aiming at re-using CO<sub>2</sub> and implementing a *validated laboratory*  
36 *technology* based on an experimental prototype apparatus (called ProGeo) producing carbon  
37 neutral methane through the chemical conversion of CO<sub>2</sub> waste flue gases using renewable  
38 energies.

39 The research is being carried out in collaboration by researchers of the laboratories of the  
40 University of Perugia and of ENEA (the Italian National Agency for New Technologies,  
41 Energy and Sustainable Economic Development of Frascati) with the support of a cluster of  
42 Small and Medium-sized Enterprises (SME)s coordinated by the Master-up srl company. In  
43 particular, the project leverages a) on the expertise of the Departments of Chemistry, Biology  
44 and Biotechnology and of Civil and Environmental Engineering of the University of Perugia  
45 on the experimental and theoretical treatment of elementary reactive and non reactive  
46 molecular processes [1-8] and b) on the engineering capabilities of designing, building and  
47 experimenting innovative apparatuses of ENEA [9] researchers with the support of the above  
48 mentioned SMEs specialized in the field of molecular science software modeling. Such

49 collaboration has been incubated by the networking activities of some COST actions [10,11]  
50 and of some virtual communities of the European Grid Infrastructure EGI.eu [12, 13].

51 The first purpose of this paper is to illustrate how the developed distributed computing  
52 machinery (i.e. the collaborative (workflowed) sharing of high level accurate and approximate  
53 ab initio and empirical calculations of the electronic structure of the involved molecular  
54 systems, the fitting and, as an alternative, force field formulations of the related potential  
55 energy surfaces together with accurate quantum, quantum-classical, quasi-classical  
56 dynamical calculations of the detailed properties of the system plus their integration with  
57 statistical averaging of the detailed properties over the unobserved parameters as well as their  
58 combination in multi-scale treatments) has provided a solid ground for a converging  
59 comparison of experimental and computational investigations of the methane formation from  
60 CO<sub>2</sub> using solid state catalyzed processes.

61 The second purpose of the paper is to illustrate how the progress made in the above  
62 mentioned research line has prompted a new one based on the use of the Synchrotron and the  
63 Free Electron Laser light sources to investigate the possibility of developing routes alternative  
64 to solid state catalyzed techniques. In particular, details are given on how CO<sub>2</sub> hydrogenation  
65 can occur via plasma generation by either electrical discharges or by vacuum ultraviolet  
66 (VUV) excitations of a CO<sub>2</sub> + H<sub>2</sub> gas mixture.

67 Accordingly, the paper is articulated as follows:

68 in section 2 structure and performances of the methanation reactor are illustrated;

69 in section 3 the computer simulation of the involved processes is analysed;

70 in section 4 the research lines driving the evolution of the apparatus towards a pure gas phase  
71 one are discussed.

72

73

74 **2. The present methanation reactor**

75 Experimental measurements of methane yields of *ProGeo* have been carried out using CO<sub>2</sub>  
76 bottles (some measurements have been performed also using CO<sub>2</sub> produced (in the ratio of 1.9  
77 kg per litre) from grapes fermentation of Marsala, a famous wine of the region where  
78 Garibaldi got ashore when starting the unification of Italy, kindly provided by IRVO (Istituto  
79 Regionale Vini e Oli)). The measurements show a complete compatibility with the catalyser  
80 used for the methanation process in *ProGeo*. The used catalyser is KATALCO<sub>JM</sub> 11-4MR a  
81 commercial product of the Johnson Matthey company (made of silicate hydrous aluminum,  
82 silicon oxide, nickel, nickel oxide, magnesium oxide, graphite porous small cylinders of  
83 average diameter 3.1 mm and height 3.6 mm containing metallic nickel) whose 2500X  
84 micrography is shown in Fig. 1.

85 The ProGeo reactor is designed for an exit CH<sub>4</sub> maximum flux of 1 Nm<sup>3</sup>/h. As shown in  
86 Figure 2, it is articulated in twin columns (twin flow channels) externally cooled thanks to  
87 shared laminar elements in order to dissipate the heat produced by the exothermicity of the  
88 reaction ( $\Delta H_{298K} = -164.9 \text{ kJ mol}^{-1}$ ). Design parameters of the reactor are:

89 GHSV (Gas Hourly Space Velocity) (ml h<sup>-1</sup> gr<sup>-1</sup>) and SV (Space Velocity) (h<sup>-1</sup>).

90 Another useful parameter is the ratio between the total flux of gas and the exposed surface of  
91 the catalyser indicated as GHCS (Gas Hourly Catalyst Surface) (cm h<sup>-1</sup>).

92 The value of the reactor parameters of the present version of the ProGeo apparatus are: GHSV  
93 = 3200 ml h<sup>-1</sup>gr<sup>-1</sup> and SV = 2100 h<sup>-1</sup> leading to a GHCS parameter of 155 cm h<sup>-1</sup>.

94 The resulting features of the present ProGeo apparatus shown in Fig. 2 are:

- 95 - External diameter of the single channel 60 mm;
- 96 - Catalysed total length 660 mm;
- 97 - Maximum inlet flux (H<sub>2</sub>+CO<sub>2</sub>) 6 Nm<sup>3</sup>/h;

- 98 - Reference Molar ratio ( $\text{CO}_2/\text{H}_2$ ) 1/5;  
99 - Maximum outlet flux ( $\text{CH}_4$ ) 1  $\text{Nm}^3/\text{h}$ .

100 Experimental measurements have been carried out by varying:

- 101 - the  $\text{CO}_2/\text{H}_2$  molar ratio from a minimum of 1/4 to a maximum of 1/5.5;  
102 - the reactor temperature from a minimum of 220 °C to a maximum of 450 °C;  
103 - the  $\text{H}_2$  flux from a minimum of 0.8  $\text{Nm}^3/\text{h}$  to a un maximum of 2.5  $\text{Nm}^3/\text{h}$ ;  
104 - the pressure from a minimum of 1 bar to a maximum of 2 bar.

105 A key element for a quantitative measurement of the reaction yield is the regular monitoring  
106 of the temperature both along the central axis of each channel (maximum temperature) and  
107 along the peripheral axis (minimum temperature). In Figure 3 the location and labels of the  
108 channel thermocouples (6 axial and 3 peripheral) are shown as dots together with the cross  
109 section of the twin columns.

110 The methanation reactions occur mainly on the top segment of the reactor close to the inlet  
111 (90% in the first 200 mm). Accordingly, most of the produced heat is detected by the central  
112 thermocouples (T5, T6 and T14, T15). The average of the measured percentage of produced  
113 methane is plotted as a function of the molar ratio and temperature in Fig. 4. As is apparent  
114 from the Figure, the threshold temperature is 240°C and the range of temperature of optimal  
115 yield is 300 - 350°C. The molar ratio has a significant impact on the percentage of produced  
116 methane with the optimum value of the  $\text{CO}_2 / \text{H}_2$  ratio being 1/5.

117

### 118 **3. Computational simulation of methane production from the $\text{CO}_2+\text{H}_2$ reaction with** 119 **solid phase catalysis**

120 In our laboratory the initial rationalization of the processes involved in producing methane via  
121 the Sabatier reaction (i.e. by making  $\text{H}_2+\text{CO}_2$  react on a catalytic surface [14]) has been

122 performed by following the flowchart of ZACROS [15], a Kinetic Monte Carlo (KMC)  
123 [16,17] software package written in Fortran 2003. ZACROS leverages on the Graph-  
124 Theoretical KMC methodology coupled with both cluster expansion Hamiltonians for the ad-  
125 layer energetics and the Brønsted-Evans-Polanyi relations for the activation energies of  
126 elementary events [16]. ZACROS enables researchers in the areas of Computational Catalysis  
127 and Surface Science to perform dynamic modelling of adsorption, desorption, surface  
128 diffusion, and reaction processes on heterogeneous catalysts. Although in our traditional  
129 GEMS (Grid Empowered Molecular Simulator) [18] approach based on distributed  
130 computing platforms the rates ( $r$ ) of elementary processes are computed by running molecular  
131 dynamics calculations based on Potential Energy Surfaces (PES)s originating from *ab initio*  
132 treatments, in the version of ZACROS implemented by us, for all the elementary steps  
133 potentially participating in the mechanism, the values of  $r$  to be used are taken from the  
134 literature and are usually given the following Transition State (TS) theory formulation [19,20]

$$135 \quad r = A \cdot \exp\left(-\frac{E_a}{k_B T}\right) \quad (1)$$

136 in which  $A$  is a pre-exponential factor (quantifying the propensity of the system to cross from  
137 TS to products) while the exponential term brings in the information on the PES of the related  
138 elementary process (either forward “f” or backward “b”) as the energy difference  $E_a$  between  
139 the stationary point of the potential Minimum Energy Path (MEP) at the transition state and  
140 that associated with the reactant asymptote with  $k_B$  being the usual Boltzmann constant and  $T$   
141 the temperature. The lowest approximation level formulates the pre-exponential factor as  
142  $k_B T/h$  where  $h$  is the Plank constant.

143 In the adsorption of gaseous species, however, the most frequently used expression for the  
144 rate of reaction is the well known Hertz-Knudsen equation:

145 
$$r_i^{\text{ad}} = S_{0,i} \cdot A_{\text{site}} \frac{p_i}{\sqrt{2\pi m_i k_B T}} \quad (2)$$

146 where  $S_{0,i}$  is the sticking coefficient (for which we take the value 1 in this work),  $A_{\text{site}}$  is the  
 147 area of the adsorption site,  $p_i$  is the partial pressure of species  $i$  and  $m_i$  its mass.

148 A more accurate formulation of the rate coefficients makes use of the partition function for  
 149 both, the intermediate state (incorporating so far the information about the remaining degrees  
 150 of freedom) and the reactants:

151 
$$A = \left( \frac{k_B T}{h} \right) \left( \frac{Q^\ddagger}{Q_r} \right) \quad (3)$$

152 where  $Q^\ddagger$  is the transition state partition function and  $Q_r$  the partition function of reactant  
 153 species.

154 Each partition function is calculated considering that rotations and translations are frustrated  
 155 (hindered), due to the fact that they are adsorbed, and therefore can be assimilated to  
 156 vibrational degrees of freedom. In that case, the vibrational partition function takes the form:

157 
$$Q_{\text{vib},X} = \prod_k \frac{\exp\left(\frac{-h\nu_k}{2k_B T}\right)}{1 - \exp\left(\frac{-h\nu_k}{k_B T}\right)} \quad (4)$$

158 where  $Q_{\text{vib},X}$  is the total vibrational partition function of the species  $X$ , and  $\nu_k$  is the frequency  
 159 of the vibrational mode  $k$ .

160 In this way ZACROS can also simulate desorption/reaction spectra at a given temperature  
 161 providing so far a rationale for designing kinetic mechanisms and understanding experimental  
 162 data.



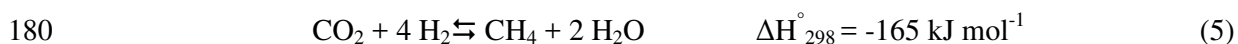
163 The elementary processes considered for the simulations are given in refs. [21,22] together  
164 with the numerical value of the related parameters. Hexagonal periodic default lattice will be  
165 the one of our choice for the Ziff, Gulari and Barshad (ZGB) reference model [23] as well as  
166 for the Sabatier Process.

167 An important outcome of our calculations is illustrated in Fig. 5 (see also ref. [19]) in which,  
168 in contrast with the suggestion of ref. [21] of a dominance of the CO\* decomposition, it is  
169 apparent from the analysis of the ZACROS calculations that the dominant process leading to  
170 the production of CH<sub>4</sub> is the hydrogenation of CO. This has suggested us to investigate  
171 possible alternative gas phase processes producing CO. For this reason, we decided to  
172 investigate the possibility of inducing CO<sub>2</sub> dissociation and producing CO+O and CO<sup>+</sup>+O<sup>+</sup>  
173 neutral and ionic chemical pairs reacting with hydrogen, respectively, via a plasma generation  
174 either by electrical discharges or by vacuum ultraviolet (VUV) photons on a CO<sub>2</sub>+H<sub>2</sub> gas  
175 mixture as will be discussed in the next section.

176

#### 177 **4. Methane production from the CO<sub>2</sub>+H<sub>2</sub> reaction without solid phase catalysis**

178 The investigations carried out on our ProGeo 20kW apparatus [22] based on the well  
179 known Sabatier reaction:



181 at moderately high pressure (2-3 atm) and high temperature (200-300°C) with the use of a  
182 solid phase catalyst (nickel, ruthenium, or alumina) has prompted us to further progress (as  
183 will be illustrated in detail in this section) to use either low cost or renewable energy to reuse  
184 waste CO<sub>2</sub> to produce methane in a circular economy scheme [22]. To this end we have  
185 undertaken the investigation of a new methanation pathway aimed at avoiding the use of the  
186 solid phase catalysis, by exploring mechanisms involving a plasma generation by electrical  
187 discharges or by vacuum ultraviolet (VUV) photons on a CO<sub>2</sub>+H<sub>2</sub> gas mixture. This effort is

188 based on the since long established expertise of our laboratories in promoting and modeling  
189 molecular processes (atomic hydrogen and excited metastable species generation by electrical  
190 discharge [24-26] to study gas phase processes induced under controlled conditions) and in  
191 fully characterizing the microscopic dynamics of elementary reactions by experimentally  
192 determining related main kinetic parameters such as rate constants, cross sections,  
193 intermolecular potentials, structure, and energy of the transition state, reaction pathways,  
194 etc. Furthermore, we can perform high resolution experiments in single collision conditions  
195 using crossed molecular beam apparatuses and studying plasma induced gas phase reactions  
196 both by microwave and RF (Radio Frequency) discharges [24-26] and by synchrotron  
197 radiation [27-29]. In particular the mentioned experiments are performed on the crossed  
198 molecular beams apparatus of the Perugia University [30,31] to measure PIES (Penning  
199 Ionization Electron Spectroscopy) data, and on the ARPES (Angle Resolved PhotoEmission  
200 Spectroscopy) end station at the GasPhase Beamline of the Elettra Synchrotron Radiation  
201 Facility (Trieste) [32,33] (more details are given later in the specific subsection).

202 The need of avoiding the use of the solid phase catalyst to perform reaction (5) and  
203 investigating possible alternative microscopic reaction mechanisms occurring in the  
204 homogeneous gas phase, is motivated by concurrency of reaction (5) and the following two  
205 main reactions:



208 as well as by the wish of avoiding unwanted reactions, like the (8), (9), (10), and (11) listed  
209 below,





214 responsible for a rapid loss of catalytic activity. Further reservations on the use of a Ni based  
215 solid phase catalyst are related to the environmental and safety issues associated with the  
216 possible formation of Ni(CO)<sub>4</sub> that is a highly toxic gaseous species produced when Ni is  
217 exposed to a gas mixture containing high pressure CO as typical of experimental conditions  
218 commonly used to maximize the yield of CH<sub>4</sub> in the methanation reaction.

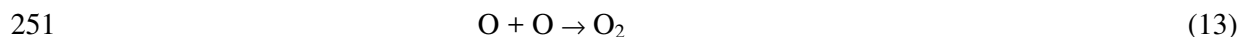
219 For this purpose we have undertaken the study of the Sabatier reaction (5) in a homogeneous  
220 gas phase environment by generating and characterizing controlled plasmas via electrical  
221 discharges and VUV photons on a pure CO<sub>2</sub> and CO<sub>2</sub>+H<sub>2</sub> gaseous mixture a technology in  
222 which our research group is leader since early 1990 [24,34,35]. In our experimental  
223 apparatuses (described in the next section) the energy pumped in the produced plasma can be  
224 controlled using both electrons and photons. In the first case we use an inhouse electrical  
225 microwave or RF discharge. In the second case we use a tunable synchrotron radiation. At the  
226 same time we can control and characterize the chemistry of the generated plasmas by studying  
227 the microscopic dynamics of the elementary chemical reactions because in molecular beam  
228 techniques [38-40] they occur in single collision regime (see again below in the specific  
229 subsection).

230

#### 231 *4.1. Experimental methods*

232 As already mentioned, a first attempt to produce plasmas containing carbon dioxide, has been  
233 made in our laboratory using a microwave discharge beam source operating in pure CO<sub>2</sub> and  
234 in an approximately 50-50% CO<sub>2</sub>-H<sub>2</sub> mixture at a global pressure of about 1600 Pa. The  
235 microwave discharge is produced in a cylindrical quartz tube (10 cm in length and 2 cm in  
236 diameter) in a brass resonance cavity (water cooled) working at 2450 MHz [24,34,35]. The  
237 applied microwave power was varied in the range 70-200 W with a reflected power lower

238 than 5%. A preliminary characterization of the produced plasma has been performed using the  
239 crossed molecular beam apparatus shown in Figure 6. Such an apparatus, usually devoted to  
240 the study of the microscopic dynamics of autoionization reactions induced by excited  
241 metastable species (also called Penning Ionization or Chemi-ionization reactions), for the  
242 experiment considered here was used by keeping the secondary beam switched off. In this  
243 way it was possible to detect only the main chemical species flowing out of the plasma  
244 microwave discharge source as an effusive molecular beam. Such an analysis was performed  
245 using the mass spectrometry characterization of the beam by means of a quadrupole mass  
246 filter located below the crossing beams region. Measured data is consistent with that  
247 previously recorded by Dobrea *et al.* [39,40] who found a large dissociation rate of carbon  
248 dioxide, according to reactions (12) and (13) below, in the case of the plasma discharge in  
249 pure CO<sub>2</sub> with respect to the one measured in CO<sub>2</sub>+H<sub>2</sub> gas mixture.



252 The CO<sub>2</sub> dissociation percentage,  $\chi$  (determined by keeping the inlet gas pressure at a constant  
253 value of 1600 Pa) was obtained by recording the CO<sub>2</sub><sup>+</sup> intensities  $I_i$  and  $I_f$  (measured with the  
254 microwave discharge off and on, respectively) using the following relationship:

$$255 \quad \chi = \frac{I_i - I_f}{I_i} \cdot 100 \quad (14)$$

256 The values of  $\chi$  amounted to about 27% in the case of the plasma discharge in pure CO<sub>2</sub>, and  
257 lowered to about 19% in the case of CO<sub>2</sub>+H<sub>2</sub> gas mixture at an applied microwave power of  
258 100 W. The measured values increased up to 52% (for plasma in pure CO<sub>2</sub>) and 36% (in  
259 plasma in CO<sub>2</sub>+H<sub>2</sub> mixture) when the microwave power was doubled to about 200 W.  
260 Measured data is reported in Table 1 and compared with that of refs. [39,40].

261 Then the same apparatus of Figure 6 was used to measure plasma-assisted CO<sub>2</sub> conversion  
262 into hydrocarbons as an alternative way of producing synthetic fuel (this technique has been

263 recently applied by Welzel *et al.* [41] who determined the experimental conditions suited to  
264 give a 20% yield for CH<sub>4</sub> formation. We better characterized the chemistry of the produced  
265 plasmas by investigating the detail of the chemical reactions induced by the ionic species that  
266 can be formed in such gaseous environments using VUV photons as a source of energy to  
267 induce excitation and ionization processes responsible for the plasma generation. The  
268 experiment was performed at the ELETTRA Synchrotron Light Laboratory (Trieste, Italy),  
269 using the ARPES end station at the Gas Phase having a very high photon intensity and a  
270 tunable wavelength. Details of the beam-line and of the end station have been already  
271 reported elsewhere [42,43] and the characteristics of the apparatus used for the experiment are  
272 discussed in detail in refs. [44,45]. Specific features of the experiment reported here are: i) the  
273 use of a synchrotron radiation tunable energy source, working in the energy range of 35-50  
274 eV with a resolution of about 1.5 meV; ii) the detection of all produced charged particles  
275 (electrons and ions) in the generated plasma, recording them in a time resolved measurements  
276 in which we are able to extract the spatial momentum components of final ions by using the  
277 electron-ion-ion coincidence technique coupled with an ion position sensitive detector (see  
278 below).

279 Figure 7 shows: i) on the left hand side panel, a scheme of the main part of the experimental  
280 apparatus based on the electron-ion-ion coincidence technique, and, as an example, the  
281 coincidence plot recorded in the double photoionization experiment of carbon dioxide at a  
282 photon energy of 44 eV (in such a panel are also shown the recorded mass spectra of the  
283 product ions at the same photon energy); ii) on the right hand side panel, a picture of such a  
284 device. As can be seen from the Figure, the monochromatic synchrotron light beam crosses at  
285 right angles an effusive molecular beam of CO<sub>2</sub> neutral precursors, and the product ions are  
286 then detected in coincidence with photoelectrons. The coincidence electron-ion-ion extraction  
287 and detection system consists in a time of flight (TOF) mass spectrometer equipped with an

288 ion position sensitive detector (stack of three micro-channel-plates with a multi-anode array  
289 arranged in 32 rows and 32 columns). As mentioned above, such a detector has been  
290 especially designed in order to properly measure the spatial momentum components of the  
291 ionic products [46].

292 Carbon dioxide was supplied from a commercial cylinder at room temperature (99.99%  
293 nominal purity) to a needle effusive beam source. An adjustable leak valve along the input gas  
294 pipe line was used in order to control the gas flow, which was monitored by checking the  
295 pressure in the main vacuum chamber.

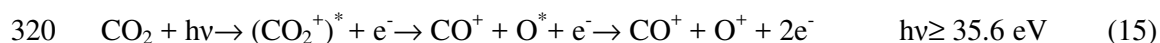
296

#### 297 *4.2. The chemical role of CO<sub>2</sub> in the generated plasma*

298         When using excitation energies higher than 35 eV, it is possible to induce ionization  
299 phenomena in gaseous mixtures containing carbon dioxide with the production of CO<sup>+</sup>, O<sup>+</sup>  
300 and CO<sub>2</sub><sup>2+</sup> ions (as a matter of fact plasmas constitute over 99% of the known matter of the  
301 Universe). Indeed, the formation of doubly charged positive ions in gas phase using VUV  
302 photons or energetic electrons as ionizing source is routine in a number of laboratory  
303 experiments able to reproduce the physical conditions of the upper atmosphere of some  
304 planets of the Solar System (such as Mars, Venus, and Titan) [27-29, 49-52]. Once these ionic  
305 species are produced in metastable states, after their typical lifetime they can, dissociate into  
306 ionic fragments having a high kinetic energy content due to the strong Coulomb repulsion  
307 characterizing the low stability of their electronic structure. This so called “Coulomb  
308 explosion” of molecular dications provides their fragments with a large amount of translation  
309 energy (several eV) that enhances the chemical reactivity of the plasma with the following  
310 consequences on the properties of a gas mixtures: i) the chemical behavior is radically  
311 changed because the removed electrons may change sensibly the electronic configuration of  
312 the neutral species and modify its chemical reactivity; ii) the interaction is much more intense

313 than the neutral-neutral one, making more likely collision events and increasing the chemical  
 314 reactivity in plasma environments; iii), the formation of fragment ions with large kinetic  
 315 energy increase chemical reactivity.

316 In particular our experiments with a plasma containing CO<sub>2</sub> confirmed the carbon dioxide  
 317 dissociation paths singled out by previous studies [53-56] and quantified the following ones  
 318 allowed to indicates that the following main processes are possible with their respective  
 319 threshold energies:



323 together with relative threshold energies. In particular, reaction (15) is an indirect process  
 324 occurring below the double photoionization threshold of 37.34 eV, via the formation an  
 325 excited state of the (CO<sub>2</sub><sup>+</sup>)<sup>\*</sup> monocation, followed by the production of an intermediate  
 326 autoionizing oxygen atom O<sup>\*</sup>. Reaction (16) produces a stable CO<sub>2</sub><sup>2+</sup> dication. Reaction (17)  
 327 mainly occurs via the formation of a long lived dication with lifetime  $\tau \geq 3.1 \mu\text{s}$ ,  
 328 corresponding to the production of CO<sub>2</sub><sup>2+</sup> ions in the ground X<sup>3</sup>Σ<sub>g</sub><sup>+</sup> electronic state with an  
 329 internal energy below the threshold towards the Coulomb explosion [53-56].

330 Figure 8 shows the relative cross sections measured in the investigated photon energy  
 331 range for reactions (15), (16), and (17) recorded using the electron-ion-ion coincidence  
 332 technique discussed in the previous section. The analysis of data, collected at each  
 333 investigated photon energy, is based on the density distribution evaluation of coincidences in  
 334 the measured coincidence spectra, such as the one reported in the left panel of Fig. 7. On such  
 335 an analysis, the total counts (recorded at each investigated photon energy) corresponding to  
 336 the ion pair peak coming out from the same reaction gives the relative cross section for the  
 337 three investigated reaction channels of Fig. 8. At the same time, the KER (Kinetic Energy

338 Released) of product ions can be evaluated by analyzing the dimension and the shape of the  
339 recorded peaks using the method suggested by Lundqvist *et al.* [57]. This procedure, applied  
340 to the  $\text{CO}^+/\text{O}^+$  coincidences signal related to reactions (15), (16) and (17), has allowed us to  
341 evaluate the KER distributions of all product ions for different values of the investigated  
342 photon energy (36.0, 39.0, 41.0, 44.0 and 49.0 eV). The results shown in Fig. 9 tell us that the  
343 KER value for the  $\text{O}^+$  ions ranges between 1.0 and 5.0 eV, while the  $\text{CO}^+$  KER can reach 3.0  
344 eV or more and, changes the maximum value depending on the investigated photon energies.  
345 These results demonstrate that the chemical reactivity of plasmas containing  $\text{CO}_2$  is strongly  
346 increased by the presence of  $\text{CO}^+$  and  $\text{O}^+$  ions having a very high kinetic energy. In particular,  
347 the fast  $\text{CO}^+$  ions, are expected to react with molecular and atomic hydrogen (both produced  
348 in a plasma generated by a microwave discharge in a gaseous  $\text{CO}_2+\text{H}_2$  mixture [24-26,  
349 39,40]) playing a pivotal role in the plasma-assisted  $\text{CO}_2$  conversion on  $\text{CH}_4$  fuel. In this  
350 respect, it has to be noted that Knott *et al.* measured reactive cross section for the  $\text{CO}^+ + \text{H}_2 \rightarrow$   
351  $\text{HCO}^+ + \text{H}$  reaction, obtaining a related rate constant value ranging between  $1.6 \times 10^{-9}$  and  
352  $3.0 \times 10^{-9} \text{ cm}^3 \text{ s}^{-1}$  in the collision energy range of 0.01 – 3.0 eV, indicating a pronounced  
353 decline at elevated collision energies, higher than 7.0 eV [58]. An analogous situation has  
354 been recorded by Farrar and coworkers in their study of  $\text{H}_2^+ + \text{CO}$  proton transfer reaction  
355 producing  $\text{HCO}^+$  in the 0.74 - 9.25 eV collision energy range [59]: at higher energies, the  
356 cross section drops rapidly whereas, at low energies, the  $\text{HCO}^+$  products are highly excited,  
357 with 90% of the available energy in internal excitation.

358         A first consideration can be made regarding the production of a low energy  $\text{CO}_2\text{-H}_2$   
359 plasma. In such a case, it has to be noted that, by using  $\text{CO}^+$  ions coming from the Coulomb  
360 explosion of  $\text{CO}_2^{2+}$  molecular dication, a projectile reactive species is available with a  
361 translational energy content of about 2.0-2.5 eV (see Fig.9). This means that in our plasma the  
362  $\text{CO}^+ + \text{H}_2$  reaction, also favored by a stronger long range trapping attraction, can occur with a



363 rate constant value higher than those of neutral-neutral reactions (having in most cases typical  
364 rate constant values of about  $10^{-10}$ - $10^{-12}$   $\text{cm}^3 \text{s}^{-1}$  [60]). In the case of the production of high  
365 energy  $\text{CO}_2$ - $\text{H}_2$  plasmas, further considerations can be made, in order to take into account that  
366 the  $\text{CO}^+$ + $\text{H}_2$  measured cross section shows a strong decrease when the collision energy  
367 becomes higher than 7 eV [58]. In this case, Knott *et al.* explain their results by invoking a  
368 dissociation process of the  $\text{HCO}^+$  product with a threshold beyond 7.0 eV. In fact, since  $\text{CO}^+$   
369 ( $^2\Sigma$  symmetry) has a predissociating state close to 8.0 eV (vertical) [61] and the vertical  
370 transition to the first excited state of  $\text{H}_2$  ( $^3\Sigma_u^+$ ) represents a step of about 9.0 eV, the  $\text{CO}^+$ + $\text{H}_2$   
371 collision should result in an excitation of  $\text{CO}^+$  rather than of  $\text{H}_2$ . We fully agree with such  
372 authors and it is important to note that the possible predissociation of  $\text{CO}^+$  appears to be  
373 energetically accessible at lower energies (about 5 eV) as demonstrated by more recent  
374 calculations by Okada and Iwata [62]. Furthermore, Nobes and Radom in 1981 [63] first  
375 calculated the energy profile for the fast isomerization reaction between  $\text{HCO}^+$  and  $\text{COH}^+$   
376 demonstrating a higher stability for  $\text{HCO}^+$  and a triangular structure of the  $[\text{HCO}]^+$   
377 intermediate complex. This corroborates the absence of  $\text{H}_2^+$ + $\text{CO}$  and  $\text{CO}+\text{H}^+$ + $\text{H}$  possible  
378 competitive products in the  $\text{CO}^+$ + $\text{H}_2$  reaction for which Knott *et al.* [58] have recorded only  
379  $\text{HCO}^+$  ions with any evidence of both  $\text{H}_2^+$  and  $\text{H}^+$  products. In their analysis such authors,  
380 following ref. [61], did not consider the possibility of the  $\text{HCO}^+$  with the rupture of the C---O  
381 bond, because their relatively low investigated collision energy range. Moreover, such a  
382 dissociation process is energetically allowed and can occur in high energy  $\text{CO}_2$ - $\text{H}_2$  plasmas,  
383 where the formed  $\text{HCO}^+$  ions are compatible with electron attachment dissociation processes  
384 towards the production of  $\text{CH}+\text{O}$  neutral reactive species, as demonstrated since 1972 by  
385 McGregor and Berry [64].

386 Finally, the efficient conversion of  $\text{CO}^+$  and  $\text{CO}_2^+$  ions into  $\text{HCO}^+$  by collisions with  
387  $\text{H}_2$  and  $\text{H}$  reactive partners (all these species can be formed in a  $\text{CO}_2$ - $\text{H}_2$  plasma), with the

388  $\text{HOC}^+ \leftrightarrow \text{COH}^+$  isomerization, has been experimentally observed by Gerlich and coworkers  
389 [65, 66] and by Tosi *et al.* [67], corroborating the observation that in plasma environments  
390  $\text{HCO}^+$  results a very stable species as it is confirmed by astronomical observations indicating  
391  $\text{HCO}^+$  as one of the most important ions in dense molecular clouds [68,69].

392 In conclusion, it appears that the production of a  $\text{CO}_2\text{-H}_2$  plasma by a microwave discharge  
393 allowed us to generate: i) a large dissociation of carbon dioxide, according to reactions (12)  
394 and (13) (see data of Table 1); and ii)  $\text{CO}_2^+$ ,  $\text{CO}_2^{2+}$ ,  $\text{CO}^+$  and  $\text{O}^+$  ionic species (see reactions  
395 (15)-(17) and Figures 8 and 9) able to react with atomic and molecular hydrogen in order to  
396 produce  $\text{HCO}^+$  ions. All such experimental evidences are compatible with the possible  
397 formation of CH and  $\text{CH}^+$  species, being the first hydrogenation step on the carbon atom, for  
398 a plasma-assisted  $\text{CO}_2$  conversion into hydrocarbons as an alternative transformation route in  
399 synthetic fuel processing.

400 Further experimental work is in progress in our laboratory in order to investigate such  
401 possibility, and to optimize the experimental conditions in an attempt to perform the  $\text{CO}_2$   
402 hydrogenation reaction (5) via an alternative microscopic mechanism with respect to the use  
403 of the solid catalyst.

404

## 405 **5. Conclusions**

406 The reported concerted efforts of University, ENEA and SMEs researchers in  
407 investigating engineering and distributed computing have shown in this paper to what extent  
408 the advances in basic research on plasmas are amenable to the assembling of an experimental  
409 prototype applying the Sabatier reaction to a homogenous gas phase catalytic environment  
410 and to the activation of a circular process turning waste  $\text{CO}_2$  flue gases into methane. The  
411 measurements performed have not only shown the viability of the proposed solution but have  
412 also allowed the extension of computer simulations to a family of innovative mechanisms.

413 Furthermore, as a result of the investigation, useful indications have been obtained on how to  
414 proceed to develop alternative solutions to the present Ni catalysed Progeo apparatus by  
415 resorting to a gas-phase-only process for the reduction of CO<sub>2</sub> to CH<sub>4</sub>. In this work new results  
416 obtained by electrical discharges into CO<sub>2</sub>-H<sub>2</sub> gas mixtures by using molecular beam  
417 technique at different pressure regimes has been presented. Such results together with VUV  
418 CO<sub>2</sub> photoionization data collected by synchrotron radiation are then of great help in  
419 identifying the microscopic mechanisms to be exploited. Finally, the production of a CO<sub>2</sub>-H<sub>2</sub>  
420 plasma by a microwave discharge allowed us to generate a large dissociation of carbon  
421 dioxide, and to produce CO<sub>2</sub><sup>+</sup>, CO<sub>2</sub><sup>2+</sup>, CO<sup>+</sup> and O<sup>+</sup> ionic species able to react with atomic and  
422 molecular hydrogen in order to produce HCO<sup>+</sup> ions. All such experimental evidences are  
423 compatible with the possible formation of CH and CH<sup>+</sup> species, representing the first  
424 hydrogenation step on the carbon atom, for a plasma-assisted CO<sub>2</sub> conversion into  
425 hydrocarbons. This is an alternative transformation route in synthetic fuel processing that we  
426 shall investigate further to the end of grounding Progeo on more robust methanation  
427 processes.

428

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641

642 **Table 1** – The percentage of CO<sub>2</sub> dissociation in microwave discharge plasma source  
 643 produced in pure CO<sub>2</sub> and in a 1:1 CO<sub>2</sub>/H<sub>2</sub> gaseous mixture for different values of applied  
 644 microwave power. The percentage values are calculated using the simple equation (2), and are  
 645 collected working at a constant pressure of 1600 Pa into the plasma source (see text). The data  
 646 are compared with those previously collected by Dobrea *et al.* in an analogous experiment  
 647 performed at a constant pressure of about 600 Pa [41,42].

648

Microwave discharge power (W)	% of CO <sub>2</sub> dissociation in pure CO <sub>2</sub> plasma		% of CO <sub>2</sub> dissociation in a 1:1 CO <sub>2</sub> /H <sub>2</sub> plasma mixture	
	This work	ref. [14]	This work	refs. [39,40]
70	19±4	---	11±4	---
100	27±3	23	19±3	14
150	40±3	---	28±3	---
200	51±2	48	36±2	33

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653 **Figure captions**

654

655 **Fig. 1.** 2500X micrography of the commercial catalyser KATALCO<sub>JM</sub> 11-4MR (by courtesy  
656 of University of Rome Tor Vergata, Rome, Italy).

657

658 **Fig. 2.** A picture of the twin columns of the ProGeo reactor.

659

660 **Fig. 3.** Cross section of the reactor twin columns and locations of the related thermocouples  
661 (red dots). On the right hand side the labels of the thermocouples.

662

663 **Fig. 4.** Average measured percentage of the methane produced in our experiments plotted as a  
664 function of the molar ratio and of the temperature with a pressure of 2 bar.

665

666 **Fig. 5.** Relative contributions of the different elementary channels to the production of  
667 methane evaluated from the minimum energy path of the PES (triangle) and from the  
668 ZACROS simulation (diamonds, squares, reverse triangles for the different temperatures  
669 connected by solid lines). The green triangle is the value suggested in ref.[24].

670

671 **Fig. 6.** The crossed molecular beam apparatus (mainly used in Penning ionization studies and  
672 adapted for plasmas generation containing carbon dioxide) in which the secondary beam was  
673 maintained off in order to detect the main chemical species coming out as an effusive  
674 molecular beam from the plasma microwave discharge source (see text).

675

676 **Fig. 7.**The electron-ion extraction and detection system used for the electron-ion-ion  
677 coincidence measurements to characterize the charged species produced in the CO<sub>2</sub> plasma by  
678 tunable synchrotron radiation. Left panel: a scheme of the set-up with a typical coincidence  
679 spectrum recorded by the double photoionization of CO<sub>2</sub> at a photon energy of 44 eV (see  
680 text). Right panel: a picture of such a device.

681

682 **Fig. 8.**Measured cross sections for the three main processes observed in the CO<sub>2</sub> plasma  
683 generation by using synchrotron radiation in the energy range of 34-50 eV and the electron-  
684 ion-ion coincidence technique (see text).

685

686 **Fig.9.**Kinetic energy distributions for CO<sup>+</sup> and O<sup>+</sup> fragment ions originating by Coulomb  
687 explosion of CO<sub>2</sub><sup>2+</sup> dication produced in the CO<sub>2</sub> plasma by tunable synchrotron radiation.

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690 FIG. 1

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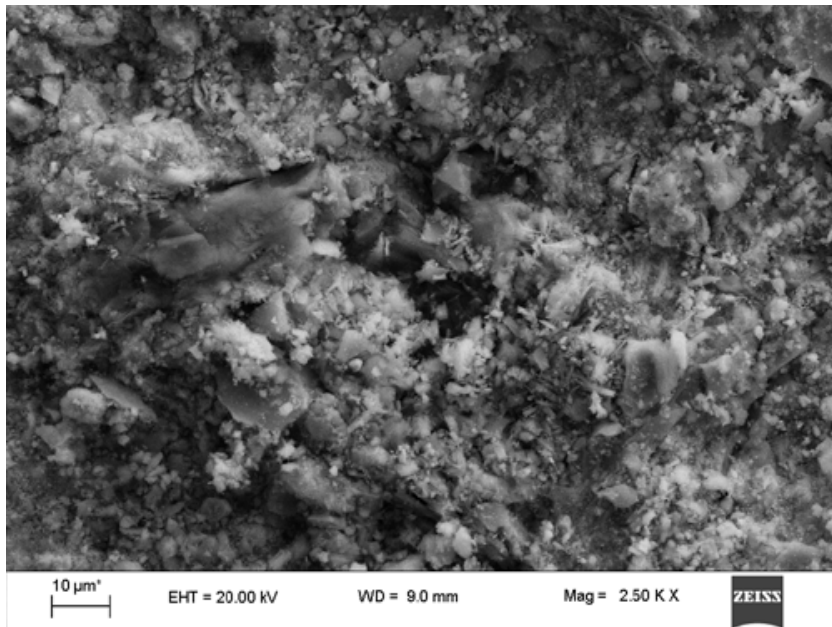
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707 FIG. 2

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731 FIG. 3

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**Methanation Reactor: inner thermocouples in channels A and B**

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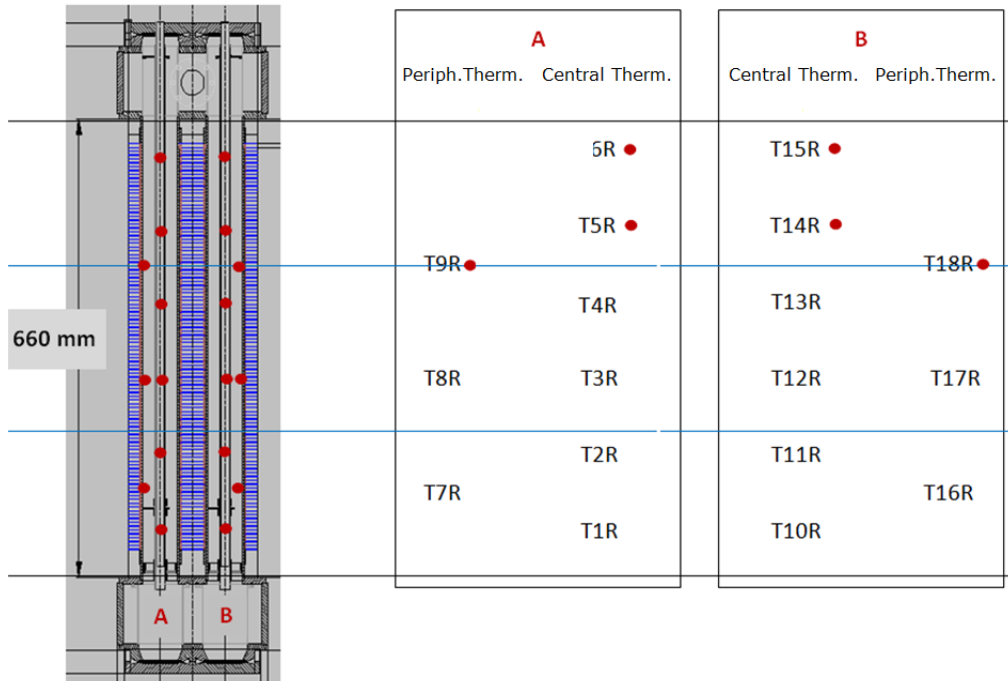
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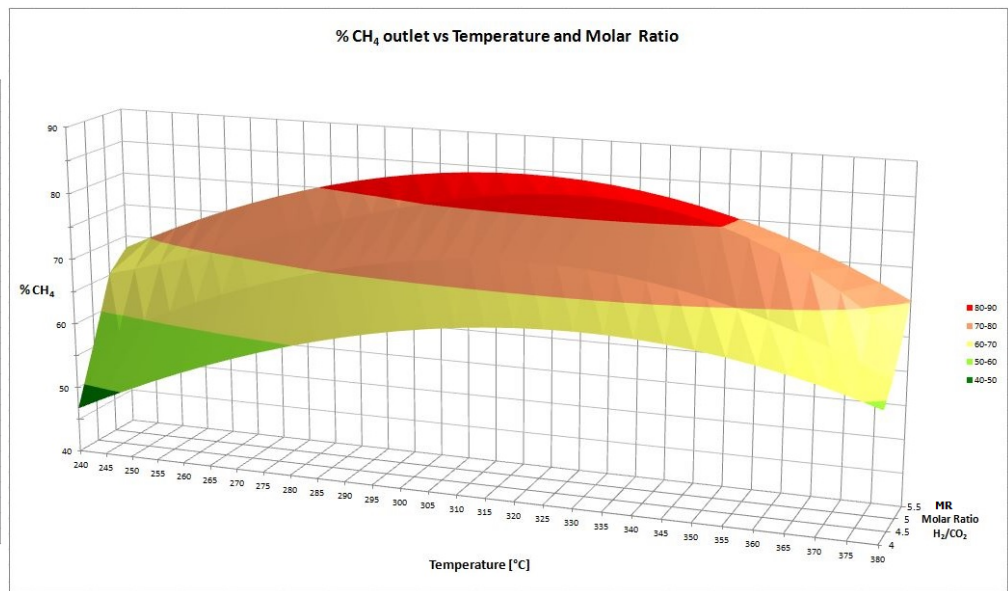
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744 FIG. 4

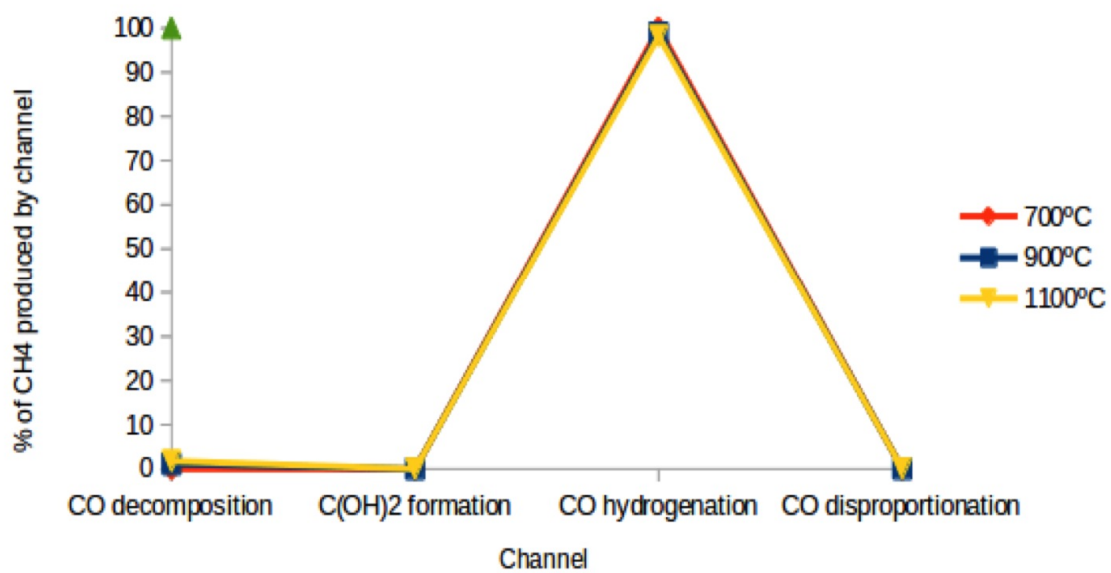
T [°C]	MR 4	4.5	5	5.5
240	47	56	65	68
245	49	58	66	70
250	51	60	68	72
255	52	62	69	74
260	54	63	70	75
265	56	65	71	77
270	57	67	73	78
275	58	68	74	79
280	60	69	75	81
285	61	70	77	82
290	62	71	78	82
295	63	72	79	83
300	64	73	80	84
305	64	73	81	84
310	65	74	82	84
315	65	74	82	84
320	66	74	83	84
325	66	74	83	84
330	66	74	83	84
335	66	73	83	83
340	66	73	82	82
345	65	72	81	82
350	65	71	80	80
355	64	70	79	79
360	64	69	77	78
365	63	67	75	76
370	62	66	72	74
375	60	64	69	72
380	59	62	66	70



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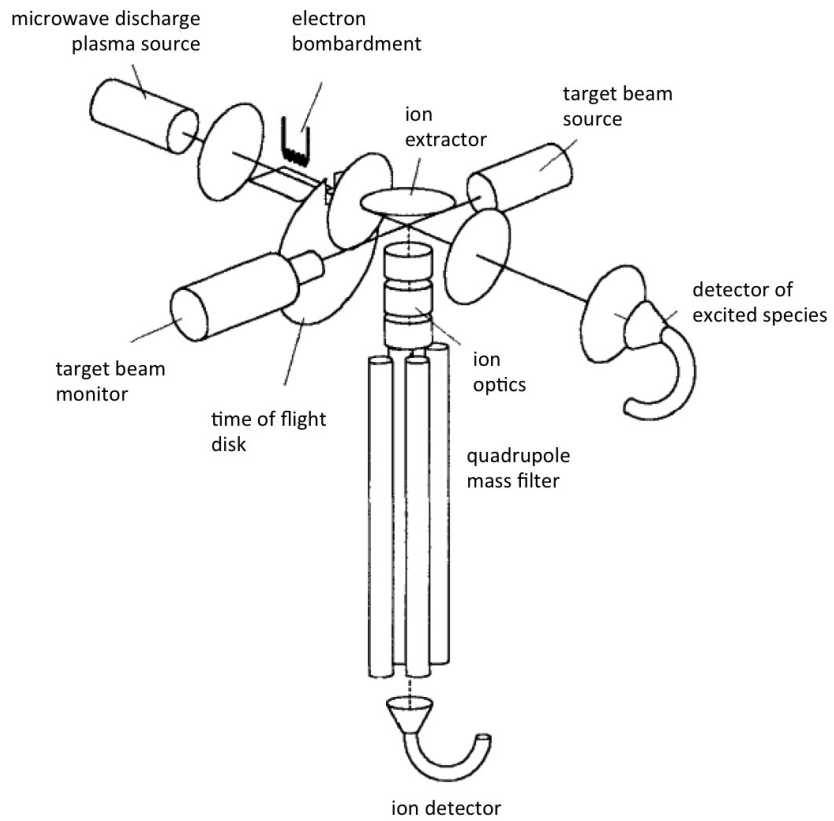
747 FIG. 5



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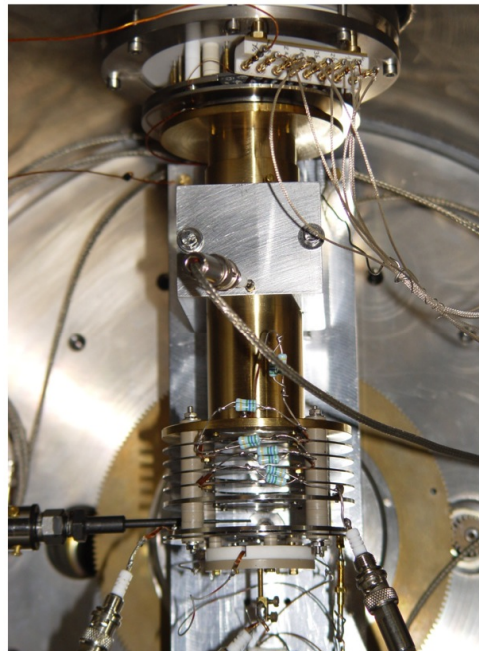
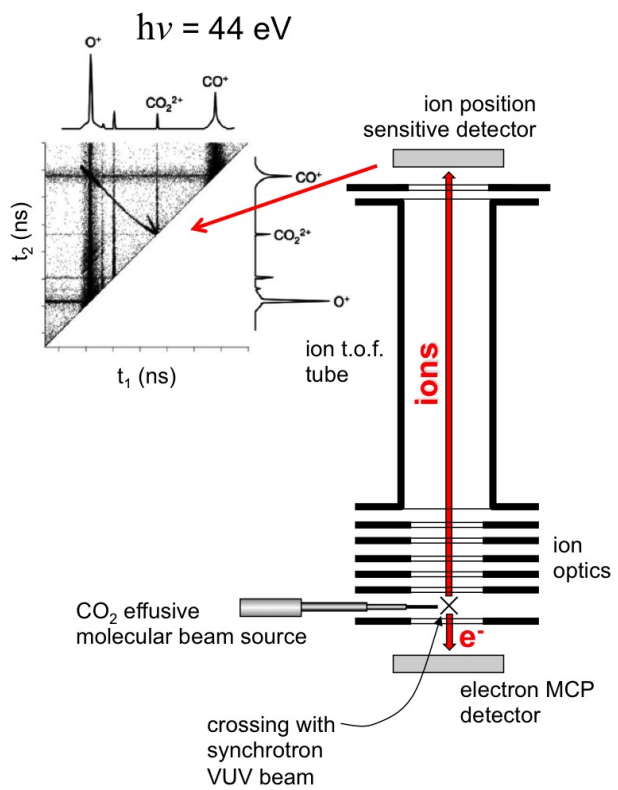
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750 FIG. 6



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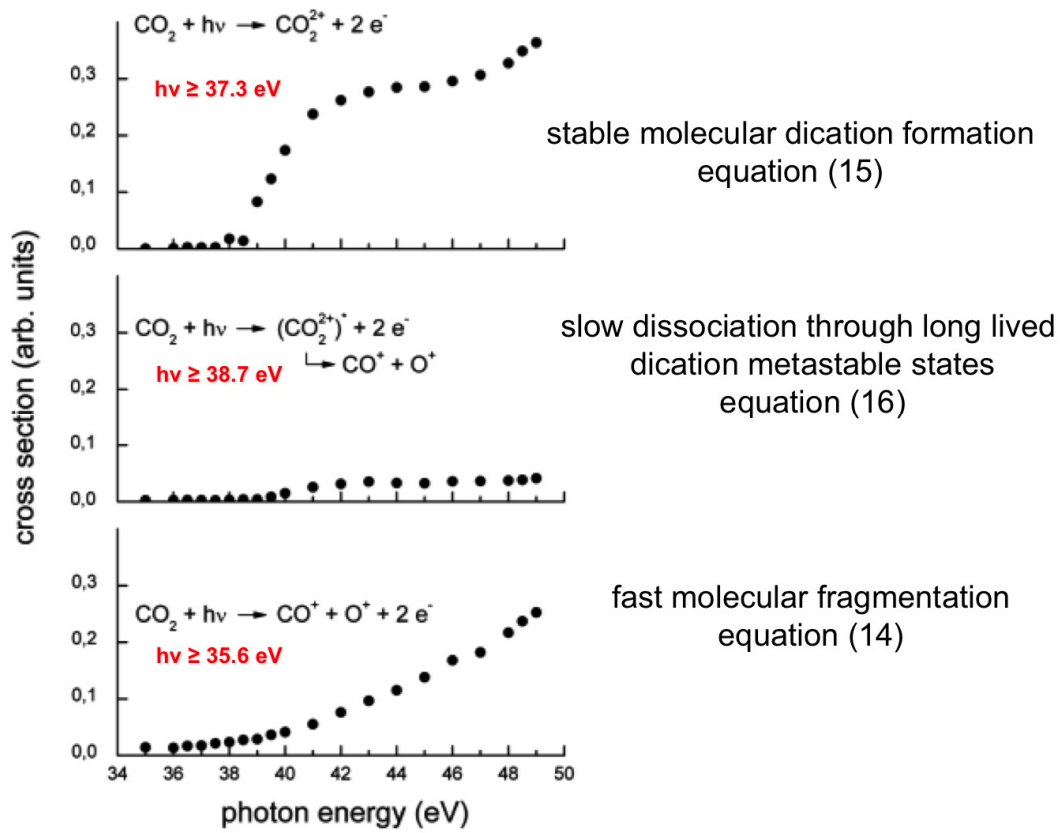
752 FIG. 7



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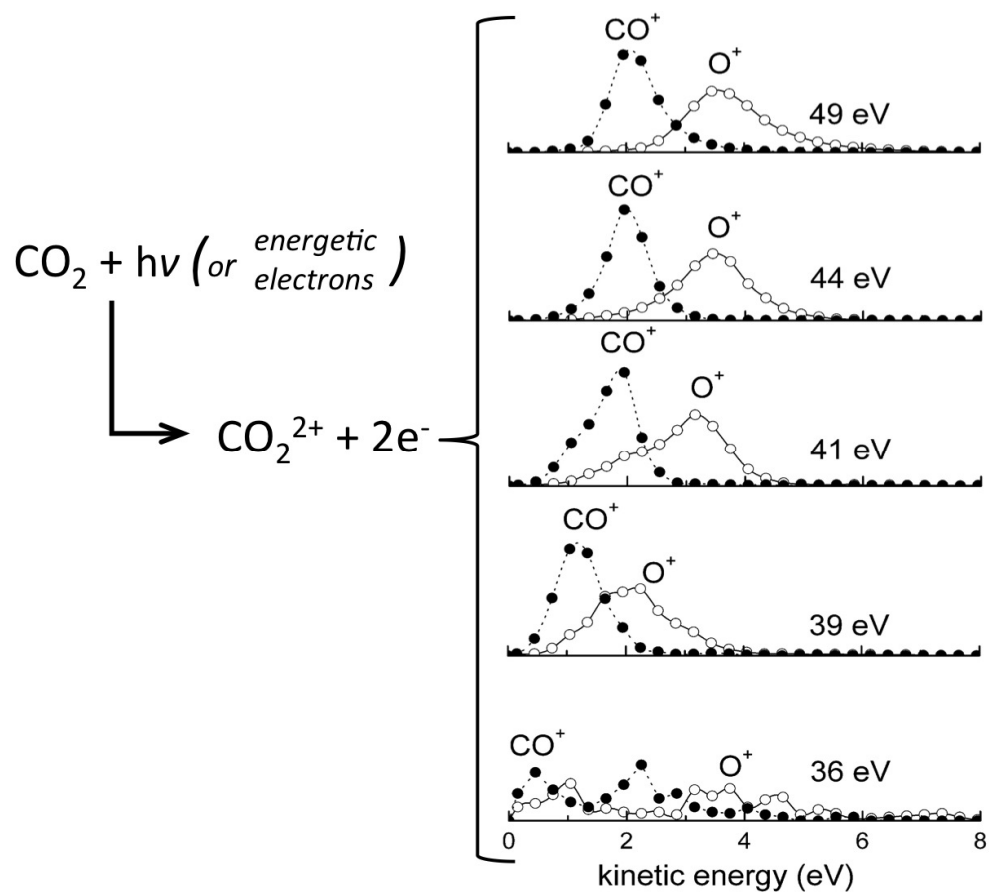
755 FIG. 8



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